## Problem set for the course "Mesoscopic Physics"

Rules: You can choose at wish from problems having the same main number, however, only one of the solutions will count (i.e., you need to choose one of these and solve only that).
Grading is as follows:

> 5: $71-$ points,
> 4: $56-70$ points,
> 3: $41-55$ points,
> 2: $31-40$ points.

## Deadline: 17. June, 5pm.

## Random Matrix Theory:

1.1 (15p) Wigner distribution for $2 \times 2$ unitary matrices: Consider a general $2 \times 2$ Hamiltonian, and parametrize it in terms of the Pauli matrices $\sigma_{0}$ (identity) and $\sigma_{x, y, z}: H=\sum_{\alpha} \Delta_{\alpha} \sigma_{\alpha}$. Construct the distance, $\mathrm{d} s^{2}=\operatorname{Tr}(\mathrm{d} H \mathrm{~d} H)$, from that the metric tensor, and show that the invariant measure is $\mathrm{d} \mu(H)=C \prod_{\alpha} \mathrm{d} \Delta_{\alpha}$. Express the eigenvalues in terms of these parameters, and also express the full Gaussian unitary distribution, $\propto \mathrm{d} \mu(H) e^{-\lambda \operatorname{Tr}\left(H^{2}\right)}$ in terms of these four real parameters. Next, introduce spherical coordinates $\left(\Delta, \theta\right.$, and $\phi$ ) for the parameters $\Delta_{x, y, z}$, compute the Jacobian, and reexpress the distribution in terms of these Gaussian variables. Integrate over the variables $\theta, \phi$, and $\Delta_{0}$, and obtain the distribution of $\Delta$. Finally, determine the average level spacing, $\delta \epsilon \equiv 2\langle\Delta\rangle$, and derive the distribution of the dimensionless level spacing, $s \equiv 2 \Delta / \delta \epsilon$.
1.2 (25p) Wigner distribution for a symplectic manyfold:
a. (10p) Assume a two-level system with two spin degenerate levels, $|a \sigma\rangle$ and $|b \sigma\rangle$. Time reversal symmetry is represented by an anti-unitary operator, which satisfies the following relations: (1) It leaves the Hamiltonian invariant, $T H T^{\dagger}=H$. (2) $T^{2}=-1$. (3) $T|a \uparrow\rangle=|a \downarrow\rangle$ and $T|a \downarrow\rangle=-|a \uparrow\rangle$. (Similar relations hold for the other state). Starting from these properties and the requirement that the Hamiltonian matrix must be Hermitian show that the most general 4 x 4 matrix describing our two-level system is:

$$
\mathbf{H}=\left(\begin{array}{cc}
\epsilon_{a} \mathbf{1} & \alpha_{0} \mathbf{1}+i \vec{\alpha} \vec{\sigma} \\
\alpha_{0} \mathbf{1}-i \vec{\alpha} \vec{\sigma} & \epsilon_{b} \mathbf{1}
\end{array}\right)
$$

with $\epsilon_{a, b}, \alpha_{0}$ and $\alpha_{x, y, z}$ real parameters. Show that the eigenvalues of this matrix are twofold degenerate, $E_{ \pm}=\left(\epsilon_{a}+\epsilon_{b}\right) / 2 \pm \Delta=\left(\epsilon_{a}+\epsilon_{b}\right) / 2 \pm \sqrt{\alpha_{0}^{2}+|\vec{\alpha}|^{2}+\left(\epsilon_{a}-\epsilon_{b}\right)^{2} / 4}$. (Hint: in spin space, use the eigenbasis of the operator $\vec{\alpha} \vec{\sigma}$.)
b. (15p)

Now proceed as in the previous exercise. Construct the distance, $\mathrm{d} s^{2}=\operatorname{Tr}(\mathrm{d} H \mathrm{~d} H)$, from that the metric tensor, and show that the invariant measure is $\mathrm{d} \mu(H)=C \mathrm{~d} \epsilon_{a} \mathrm{~d} \epsilon_{b} \prod_{k} \mathrm{~d} \alpha_{k}$. Express the full Gaussian unitary distribution, $\propto \mathrm{d} \mu(H) e^{-\lambda \operatorname{Tr}\left(H^{2}\right)}$ in terms of these six real parameters. Next, introduce spherical coordinates $(\Delta$ for the splitting, and $\theta_{1}, \theta_{2} \theta_{3}$ and $\phi$ for the $\alpha$ 's and $\left.\left(\epsilon_{a}-\epsilon_{b}\right) / 2\right)$. Compute the Jacobian, and reexpress the distribution in terms of these variables. Integrate over the angles, and obtain the distribution of $\Delta$. Finally, determine the average level spacing, $\delta \epsilon \equiv 2\langle\Delta\rangle$, and derive the distribution of the dimensionless level spacing, $s \equiv 2 \Delta / \delta \epsilon$.
$2.1(\mathbf{1 5 p}+\mathbf{1 0 p})$ You can use Mathematica or MatLab for this exercise, or any other programming language. If you code in C or Pascal, you could consider using the numerical recipes routines for diagonalization. Consider the following two dimensional Hamiltonian with periodic boundary conditions:

$$
H=-\sum_{i=1}^{N} \sum_{j=1}^{N}\left(c_{i, j}^{\dagger} c_{i+1, j}+c_{i, j}^{\dagger} c_{i, j+1}+h . c .\right)+\sum_{i, j} \epsilon_{i, j} c_{i, j}^{\dagger} c_{i, j}
$$

where the $\epsilon_{i, j}$ 's denote independent random variables in the interval $[-1,1]$.
a. (8p) Diagonalize numerically the above Hamiltonian for various realizations of the disorder. [Construct operators $\Phi_{E} \equiv \sum_{i} \phi_{i}(E) c_{i}$ that satisfy $\left.\left[H, \Phi_{E}\right]=-E \Phi_{E}\right]$. To have reasonable run times, use values $N \sim 30$.
b. (7p) Do the statistics for the level spacings $s$ for energies in the interval $E \in[-0.2,0.2]$. (For a given realization of disorder, compute $s_{i}=E_{i+1}-E_{i}$ for $E_{i} \in[-0.2,0.2]$ (allow $E_{i+1}$ to be outside this interval, otherwise you introduce artificial cut-off errors). Do that for several realizations of the disorder, until you get good statistics. Be careful with the normalization of the numerically computed distribution function. For instance, if you use a mesh of spacing $\Delta s$ to do the statistics, then you have $P\left(s \in\left[s_{0}-\Delta s / 2, s_{0}+\Delta s / 2\right]\right) \approx P\left(s_{0}\right) \Delta s$, with $P\left(s_{0}\right)$ the probability density you want to estimate at point $s_{0}$. Make sure that $\Delta s$ is small enough (must be smaller than typical level spacing, $\delta \epsilon!$ ).
c. $(+\mathbf{1 0 p})$ Determine the universal correlation function $\left.C\left(s-s^{\prime}\right):\left\langle\varrho(E) \varrho\left(E^{\prime}\right)\right\rangle=\langle\varrho(E)\rangle\left\langle\varrho\left(E^{\prime}\right)\right\rangle C\left(\left(E-E^{\prime}\right) / \delta \epsilon\right)\right)$.
$2.2(10 p)$ Consider free electrons $(m=\hbar=1)$ in a two-dimensional box of size $L_{x}=1000$ and $L_{y}=1000 \pi$. Make a statistics of the level spacing for states with energy $1<E<1.1$. Measure energy separations in the average energy separation $\Delta, s \equiv \Delta E / \Delta$. Estimate $\Delta$ analytically and compare it to your numerics. Discuss the distribution function $P(s)$.
[Hint: You will have to create a vector variable that will contain the energies of states with energies $1<E<1.1$. Then you have to order the energy of these states. You will have to create another array to make the statistics, where you just compute, how many states you have with separation $s$. Be careful, and properly normalize the distribution function $P(s)$. You can solve this problem in any programming language you like.]

## 3. (20p) Weak localization correction for a cavity with time reversal symmetry

In this problem, we shall determine the weak localization correction to the average conductance of a cavity within the circular orthogonal ensemble (COE).
a. (6p) First, analyze the structure of the S-matrix. Repeat the derivation at class to show that, in the presence of time reversal symmetry (and in the absence of spin), the S-matrix of a cavity is symmetrical, $S=S^{T}$. Show that any unitary matrix can be represented as $S=U D U^{\dagger}$, with $U$ a unitary matrix, and $D$ a diagonal matrix containing the eigenvalues of $S$, all on the unit circle, $s_{\alpha}=e^{i \phi_{\alpha}}$. Now prove that - with an appropriate phase choice - any symmetrical unitary matrix can be written as $\Omega^{T} \Omega$, with $\Omega$ a unitary matrix. Argue that $\Omega \in \operatorname{CUE}$.
b. (2p) Now follow the derivation at class, and consider a cavity with $N$ channels on the left and $N$ channels on the right. Using the Landauer-Büttiker formula, $g=h G / e^{2}=T=\sum_{i \in L} \sum_{j \in R}\left|S_{i j}\right|^{2}$, show that the expectation value of the transmission is given as

$$
\left.\langle T\rangle=2 N^{3} M, \quad \text { with } M=\left.\langle | \Omega_{\alpha i}\right|^{2}\left|\Omega_{\alpha j}\right|^{2}\right\rangle_{i \neq j}
$$

c. (12p) Determine $M$. Use the fact, that every column of $\Omega$ can be thought of as a real unit vector of dimension $d=4 N, x=\left\{\operatorname{Re} \Omega_{11}, \operatorname{Im} \Omega_{11}, \operatorname{Re} \Omega_{12}, \ldots\right\}=\left\{x_{1}, x_{2}, \ldots, x_{4 N}\right\}$. Introduce the surface of the $d$-dimensional unit sphere as

$$
A_{d} \equiv \int \mathrm{~d} x_{1} \ldots \mathrm{~d} x_{d} \delta\left(1-\left(x_{1}^{2}+\cdots+x_{d}^{2}\right)\right)
$$

Introduce furthermore the integrals, $I_{n} \equiv \int_{0}^{\pi} \mathrm{d} \theta \sin ^{n}(\theta)$. Show that the following relations hold:

$$
I_{n}=\frac{n-1}{n} I_{n-2}, \quad A_{n}=I_{n-2} A_{n-1} .
$$

Using these relations evaluate $\left.\left.\langle | \Omega_{11}\right|^{2}\right\rangle$ and show that it is simply $1 / 2 N$. Then, to evaluate $M$, first show that $\left.M=\left.\langle | \Omega_{11}\right|^{2}\left|\Omega_{12}\right|^{2}\right\rangle=4\left\langle x_{1}^{2} x_{2}^{2}\right\rangle$. Next show using the fact that $x$ is a unit vector that

$$
d\left\langle x_{1}^{4}\right\rangle+d(d-1)\left\langle x_{1}^{2} x_{2}^{2}\right\rangle=1
$$

Evaluate then $\left\langle x_{1}^{4}\right\rangle$ and show that it is $\left\langle x_{1}^{4}\right\rangle=3 /(d(d+2))$, from which one obtains

$$
M=\frac{1}{2 N(2 N+1)}, \quad \text { and } \quad\langle T\rangle=\frac{N}{2} \frac{2 N}{2 N+1}
$$

This formula implies that, in the presence of time reversal symmetry, particles entering the chaotic cavity preferentially leave it towards the lead they came from. [Hint: evaluate for $d=4 N$

$$
\left\langle x_{1}^{2}\right\rangle=\frac{1}{A_{d}} \int \mathrm{~d} x_{1} \ldots \mathrm{~d} x_{d} x_{1}^{2} \delta\left(1-\left(x_{1}^{2}+\cdots+x_{d}^{2}\right)\right)
$$

by observing that the first term, $1-x_{1}^{2}$, can be removed from the Dirac delta by appropriately rescaling $x_{2}, \ldots, x_{d}$, and then change variables, $x_{1} \rightarrow \cos (\theta)$. ]

## Transport through quantum dots

4. (15p) Büttiker's formula for frequency dependent current noise. For a single channel, evaluate the frequency dependent noise, $S^{>}(\omega)$, and plot it for (a) $V=0$ bias at $T \neq 0$ and (b) for $T=0$ and $V \neq 0$ as a function of frequency. (Use the general expression we derived at class for $\left\langle I(t) I\left(t^{\prime}\right)\right\rangle$ and take its Fourier transform. The remaining integrals can be evaluated analytically.) Discuss the graphs you get!
5.1 (20p) Rate equation for a quantum dot (grain) with dense levels.

Repeat the derivation of the master equation discussed at class for the SET for the case where you consider the three possible states, $\hat{N}=-1,0,1$. Introduce the corresponding probabilities, $P_{0}$ and $P_{ \pm}$. For the sake of simplicity assume that the SET is completely symmetrical. Denote the difference of the charging energies by $\Delta E_{ \pm} \equiv E_{ \pm}-E_{0}$.
a. (10p) Construct the following steady state master equations:

$$
\begin{align*}
& \frac{d P_{+}}{d t}=P_{0} W_{0 \rightarrow+}-P_{+} W_{+\rightarrow 0}  \tag{1}\\
& \equiv 0  \tag{2}\\
& \frac{d P_{0}}{d t}=P_{+} W_{+\rightarrow 0}+P_{-} W_{-\rightarrow 0}- P_{0} W_{0 \rightarrow+}-P_{0} W_{0 \rightarrow-}  \tag{3}\\
& \equiv 0 \\
& \frac{d P_{-}}{d t}=P_{0} W_{0 \rightarrow-}-P_{-} W_{-\rightarrow 0}
\end{align*} \overline{ } \equiv \text {. }
$$

Determine the rates $W$ as we did at class (assuming a continuum of states, and a corresponding dimensionless conductance, $g_{L}=g_{R}=g$ ). Then determine the stationary values of $P_{0, \pm}$ from the above equations. [Note that at class we defined the transition probabilities slightly differently, and they contained the probabilities $P_{ \pm, 0}$.] Remember that the rates have two contributions, corresponding to tunneling processes from/to the right and left, and introduce accordingly $W_{\alpha \rightarrow \beta}^{L}$ and $W_{\alpha \rightarrow \beta}^{R}$. This is useful for the second part of this problem.
b. (10p) Now compute and plot the current through the device as a function of bias voltage $V$ for various values of the dimensionless gate voltage $n_{G}$. [Use that $E_{N}=E_{C}\left(N-n_{G}\right)^{2}$.] Plot the linear conductance as a function of $n_{g}$, and finally the differential conductance as a function of $V$ for various values of $n_{G}$. Show that you indeed have a diamond structure as argued at class. Check the energies emerging in the differential conductance. [For these plots you can use $T=0.1 E_{C}$, e.g..]
$\mathbf{5 . 2}(\mathbf{2 1}+\mathbf{9 p})$ Rate equation for a quantum dot with a single level (simpler that 3.1 !).
Consider a quantum dot described by the Hamiltonian $H=H_{\text {lead }}+H_{\text {dot }}+H_{t}$ :

$$
\begin{align*}
& H_{\text {lead }}=\sum_{\xi, \sigma}\left(\xi+\mu_{L}\right) L_{\xi, \sigma}^{\dagger} L_{\xi, \sigma}+" L \leftrightarrow R^{\prime \prime}  \tag{4}\\
& H_{\text {dot }}=\sum_{\sigma} E_{+}|\sigma\rangle\langle\sigma|+E_{0}|0\rangle\langle 0|  \tag{5}\\
& H_{t}=t_{L} \sum_{\xi, \sigma}\left(L_{\xi, \sigma}^{\dagger}|0\rangle\langle\sigma|+\text { h.c. }\right)+" L \leftrightarrow R^{\prime \prime} . \tag{6}
\end{align*}
$$

Here $|\sigma\rangle$ and $|0\rangle$ denote the dot states with one electron of $\operatorname{spin} \sigma$ and no electrons, respectively. Note that we assumed a large level spapcing in this case, so this is just the opposite of the limit we considered at class ! Now repeat the derivation of the master equation for this system.
a. ( $\mathbf{7 p}$ ) Show that the rate of an electron jumping into the dot from the left is given by

$$
W_{0 \rightarrow+}=2 \Gamma_{L} f\left(\Delta E-\mu_{L}\right) P_{0}
$$

in the Born approximation, with $\Gamma_{L}=2 \pi \varrho_{0} t_{L}^{2}$ the width of the level, $\Delta E=E_{+}-E_{0}, f$ the Fermi function, and $P_{0}$ the probability that the dot has no electron. Note the prefactor in front, that is due to the spin! Write down the equations for all other electron transfer processes.
b. ( $\mathbf{7 p}$ ) Construct the master equation for the probabilities $P_{0}$ and $P_{+}$of having no or one electron on the dot. Determine the stacionary solution and express the current from these equations.
c. $(\mathbf{7 p})$ Now determine the linear conductance of the dot as a function of temperature and $\Delta E$, show that the Coulomb blockade peak is asymmetrically shifted. What is it's height? Discuss the problem with the peak height's temperature dependence! What do you think the solution is?
d. $(+\mathbf{9 p})$ Optional: Generalize the calculation for the case where the energies of the spin up and spin down dot states are split by a magnetic field, $E_{\uparrow}-E_{\downarrow}=B$. Compute the polarization of the current. Show that the current is spin polarized if $B>T$.
6.1 (10p) Compute the energy of a single electron transistor shown in the figure as a function of gate and bias voltages, and for a fixed number $n$ of electrons on the island. [Hint: Assume first that no charge has been transferred through the capacitor $C_{R}$, and that tunneling only occurred through $C_{L}$. Then discuss what happens if some electrons DO tunnel through this capacitor.]


## 6.2 (15p)

Consider two identical unbiased SET's in series. (See figure.) Assume that all capacitors are the same.

a. (5p) Derive the energy of the circuit for fixed number of particles $n_{1}$ and $n_{2}$ on the two dots.
b. (5p) Make a plot on the $V_{1}, V_{2}$ plane, and draw the regions of the various minimum energy states (indicate in every regime the optimum values of $n_{1}$ and $n_{2}$.
c. (5p) Can you make a turnstyle of the double dot system in series, i.e. a transistor where you transfer charge from one side to the other by applying an AC modulation on $n_{g 1}$ and $n_{g 2}$ in every cycle, so that the current is just $I=e \omega$ ?

