# Simulations in Statistical Physics Course for MSc physics students

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## Molecular dynamics

#### Program:

- ► Have an algorithm to calculate forces
- Get list of interacting particles
- Determine accelerations and velocities; step particles
- Set temperature

## Temperature

#### Definition:

- Encyclopedia Britannica, Wikipedia:
  "A temperature is a numerical measure of hot or cold."
- Thermodynamics:Second law of thermodynamics & Carnot engine

$$\delta Q = extit{TdS}$$
  $\eta_{\sf max} = \eta_{\sf Carnot} = 1 - extit{T}_C/ extit{T}_H$ 

Statistical physics:

$$\beta \equiv \frac{1}{k_B} \left( \frac{\partial S}{\partial E} \right)_{V,N} = \frac{1}{k_B T}$$



## Definition of temperature

Temperature is a measure of the random submicroscopic motions and vibrations of the particle constituents of matter.

The average kinetic energy per particle degrees of freedom is

$$\bar{E} = \frac{1}{2} k_B T$$

Molecular dynamics conserves only the total energy!

Task: Control kinetic energy!

## Setting temperature

- Experiment
  - Environment
  - lacktriangleright Mixing ightarrow uniform temperature
- Simulation
  - Control the kinetic energy (velocities)
  - $\blacktriangleright \ \, \mathsf{Mixing} \to \mathsf{Maxwell}\text{-}\mathsf{Boltzmann} \,\, \mathsf{distribution}$

#### Nosé-Hoover thermostat

Original Hamiltonian

$$H_0 = \sum_i \frac{\mathbf{p}_i^2}{2m_i} + U(\mathbf{q})$$

Heatbath in the Hamiltonian:

$$H_n = \sum_i \frac{\mathbf{p}'_i^2}{2m_i} + U(\mathbf{q}') + \frac{p_s^2}{2Q} + gk_B T \log(s)$$

- Extra degree of freedom s.
- $lackbox{ }Q$  "mass" related to s o controls the speed of convergence
- g = 3N the number degrees of freedom
- p' and q' are virtual coordinates

#### Nosé-Hoover thermostat

▶ Virtual coordinates, vs. original ones:

$$egin{aligned} \mathbf{p} &= \mathbf{p}'/s \ \mathbf{q} &= \mathbf{q}' \ t &= \int rac{1}{s} dt' \end{aligned}$$

► Solution of the new Hamiltonian:

$$\xi = \dot{s}/s = p_s/Q$$

$$\dot{\mathbf{q}}' = \frac{\mathbf{p}'}{m}$$

$$\dot{\mathbf{p}}'_i = -\frac{\partial U}{\partial q_i'} - \xi \mathbf{p'}_i$$

$$\dot{\xi} = \frac{1}{Q} \left( \sum_i \frac{\dot{\mathbf{p}}'_i^2}{m_i} - g k_B T \right)$$

## Molecular dynamics

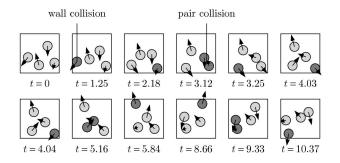
- Create sample
  - Crystal
  - Random deposition
  - Distorted crystal
  - Simulation
- Temperate sample
- Make test
- Collect data
  - ▶ Data size: e.g.  $N = 10^4$ ,  $t = 10^6$  small simulation:
  - ▶ 1 hour on 1 core PC
  - 3 doubles/atom → 24 bytes/atom/timesteps
  - Result 2.4  $10^{11}$  bytes = 240 Gigabytes

#### **Alternatives**

- Event Driven Dynamics
- Contact Dynamics
- Kinetic Monte Carlo

## Event driven dynamics

- Hard core interactions
- Interactions short in time compared to flight
- (MD needs  $\sim 20-50$  timesteps per collision, overlap of  $10^{-3}d$ )
- ▶ Integrable path  $\rightarrow$  do it



## Event driven algorithm

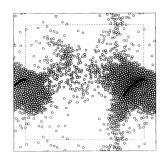
- No gravity
- ▶ Particles:  $\mathbf{r}_i(t)$ ,  $\mathbf{v}_i(t)$ ,  $\omega_i(t)$ ,  $R_i$
- ▶ Calculate collision time: Let  $\mathbf{d}_{ij} = |\mathbf{r}_i \mathbf{r}_j| R_i R_j$ , Then

$$\tau_{ij} = \frac{|\mathbf{d}_{ij}|^2}{(\mathbf{v}_i - \mathbf{v}_j)\mathbf{d}_{ij}}$$

- lacktriangle Order collision times, get the smallest  $au_c = \min_{ij} ( au_{ij})$
- Go to time  $t + \tau_c \mathbf{r}_i(t + \tau_c)$
- ullet Calculate velocities after collision  $oldsymbol{v}_i(t+ au_c)$  (may be hard...)
- Restart loop
- Next time Calculate collision time only with i, j
- Dynamic list, change only newly calculate collision times

## Inelastic collapse

- ▶ Coefficient of restitution  $r = v_n(t_c+)/v_n(t_c-)$
- Energy is lost in an exponential way (Ping Pong)
- Infinite collisions in finite time
- ▶ Solution  $\rightarrow$  r=1 if collisions occur more frequently than a parameter  $t_{\rm cont}$ , the contact duration
- Contact → small vibration :-(well...)

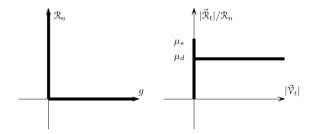


## Contact dynamics

- Perfectly rigid particles
- Constraints
- Implicit forces

$$\mathbf{v}_i(t + \Delta t) = \mathbf{v}_i(t) + \frac{1}{m_i} \mathbf{F}_i(t + \Delta t) \Delta t$$

$$\mathbf{x}_i(t + \Delta t) = \mathbf{x}_i(t) + \mathbf{v}_i(t + \Delta t) \Delta t$$



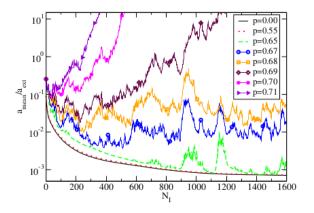
## Contact dynamics, force calculation

▶ Two particles with gap g

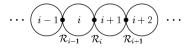
$$\begin{split} & \underbrace{\vec{F}_{l}^{\text{ret}}}_{l} \underbrace{1} \underbrace{\vec{F}_{l}^{\text{ext}}}_{l} \underbrace{\vec{F}_{l}^{\text{ext}}}_{l} \\ & \text{if} \quad \mathcal{V}^{\text{free}}_{n} \Delta t + g^{\text{pos}} > 0 \\ & \text{then} \quad \left\{ \vec{\mathcal{R}}^{\text{new}} = 0 & \text{(no contact)} \\ & \\ & \mathcal{R}^{\text{new}}_{n} = -\frac{1}{\Delta t} m_{n} \left( \frac{g^{\text{pos}}}{\Delta t} + \mathcal{V}^{\text{free}}_{n} \right) \\ & \left| \vec{\mathcal{R}}^{\text{new}}_{l} \right| > \mu \mathcal{R}^{\text{new}}_{n} \\ & \text{then} \quad \left\{ \vec{\mathcal{R}}^{\text{new}}_{l} = \mu \mathcal{R}^{\text{new}}_{n} \frac{\vec{\mathcal{R}}^{\text{new}}_{l}}{\left| \vec{\mathcal{R}}^{\text{new}}_{l} \right|} \right. \end{aligned} \tag{sticking contact)} \end{split}$$

#### Iterative solver

- ► Updates:
  - Parallel: calculate all contacts with old values then change to new at once → serious instabilities
  - Serial: update contacts one-by-one in random order



#### Particle chain



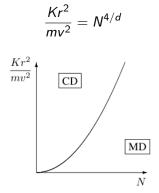
One iteration step:

$$\mathcal{R}_i^{\text{new}} = \frac{1}{2} \left( \mathcal{R}_{i-1}^{\text{new}} + \mathcal{R}_{i+1}^{\text{new}} \right),$$

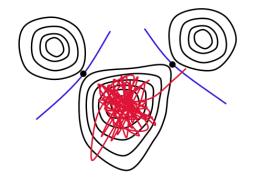
- Discretized one-dimensional diffusion equation
- lacktriangle Model of rigid particles ightarrow elastic
- Elasticity depends on the number of iterations

# Molecular versus Contact dynamics

Limit



- ▶ Particle sits in a potential well for ages . . .
- ▶ What to do?



- Long lasting steady state positions
- Slow thermally activated processes
- Infrequent-event system

#### Solution:

- Consider only jumps between neighboring energy wells
- Probability of jump  $P \sim \exp(-\beta E_b)$
- ▶ Rate of jump  $i \rightarrow j$ ,  $k_{ij} = E_b$ .



- Get all possible rates k<sub>i</sub>
- ▶ Calculate the cumulative function  $K = \sum_i k_i$
- ► Get a uniform random number *u* (between 0 and 1)
- ► Execute the event i for which  $\sum_{j=1}^{i} k_j > u > \sum_{j=1}^{i-1} k_j$
- ▶ Get new uniform random number u' (between 0 and 1)
- ▶ Update time to  $t = t + \Delta t$ ,  $\Delta t = -\log(u')/k_i$
- Recalculate rates, which have changed
- Restart loop

#### Monte Carlo

Why Monte Carlo? → Random numbers play big role!



- Rates
  - Physics
  - Molecular dynamics
- Must include all rates!
- Used for:
  - Surface diffusion
  - Surface growth
  - Syntering
  - Domain evolution

Example....

#### Methods

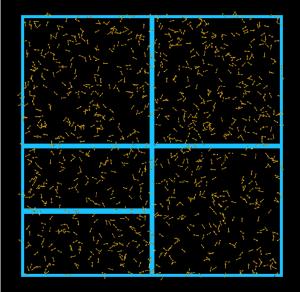
- Molecular Dynamics
  - General
- Event Driven Dynamics
  - Hard objects, at low density
- Contact Dynamics
  - Rigid particles
- Kinetic Monte Carlo
  - Infrequent events, bonded particles

#### Parallelization

- ► Why?
  - ▶ The speed of one core processor is limited
  - Larger system sizes
  - Multi-core processors
  - On multi-core system inter-processor data change is fast
- ▶ Why not?
  - Computing power is lost
  - Much more code development
  - Very often ensemble average is needed
  - Inter-computer communication is terribly slow

RAM  $ightarrow \sim 15$ GB/s, Ethernet 125MB/s, Infiniband  $\sim 1$ GB/s

# Parallelization (Bird flocking model)

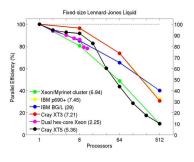


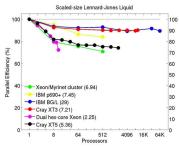
#### Parallelization

- Molecular dynamics
  - Short range interactions: Box must be duplicated, Verlet in parallel
  - Long range: Parallel fast Fourier transformation
- Contact dynamics
  - Short range interactions: Box must be duplicated, Iteration in parallel
- Event Driven Dynamics
  - List must be global, no way!
- Kinetic Monte Carlo
  - List must be global, no way!

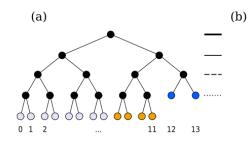
## Efficiency of parallelization

- Large systems are needed
- Boundary must be minimal





## Efficiency of parallelization



5	7	12	13
4	6	0	11
1	3	9	11
0	2	8	10

- Calculate time spent in a branch
- Calculate  $\sigma_T = \sqrt{\langle T^2 \rangle \langle T \rangle^2} / \langle T \rangle$
- ▶ Move line if necessary  $(\sigma_T > \sigma_T^*)$
- ▶ Lower in tree (up in Fig), larger the mass of the border
- Only rarely, data transfer is expensive